

1. A particle is in a state $|\psi\rangle$ when the observable corresponding to $\hat{\Omega}$ is measured. Show that the following two statements are equivalent:

(I) the probability of getting a result ω_i is $|\langle\omega_i|\psi\rangle|^2$;

(II) the expectation value (ensemble average) is $\langle\psi|\hat{\Omega}|\psi\rangle$.

[Do NOT assume that $|\psi\rangle$ is an eigenstate of $\hat{\Omega}$ – the equivalence is trivial in that case.]

2. At a given time, a particle is in a state $|\phi_0\rangle$, with

$$\phi_0(x) = \left(\frac{1}{\sqrt[4]{\pi a^2}}\right) e^{-x^2/2a^2},$$

and $a = \sqrt{\hbar/m\omega}$. A measurement is made of the momentum.

What is the probability of getting an answer within a small range $\delta p = \hbar/100a$ centred on the value \hbar/a ?

3. A hydrogen atom is prepared in an initial state described by the wave function

$$\psi_I(\mathbf{r}) = \frac{1}{\sqrt{96\pi a_0^5}} r \exp\left(-\frac{r}{2a_0}\right),$$

where a_0 is the Bohr radius. A measurement is made of the energy.

What is the probability of obtaining the ground state energy, -13.6 eV?

What other values of the energy might be obtained?

What would the probability be if instead the initial state were $\sqrt{3} \cos(\theta)\psi_I(\mathbf{r})$?

[Hint: You will need the ground-state wave function of hydrogen to answer this; it is proportional to $\exp(-r/a_0)$. The first answer works out as about 0.23.]

4. In a particular orthonormal basis, the Hamiltonian of a certain system is represented by

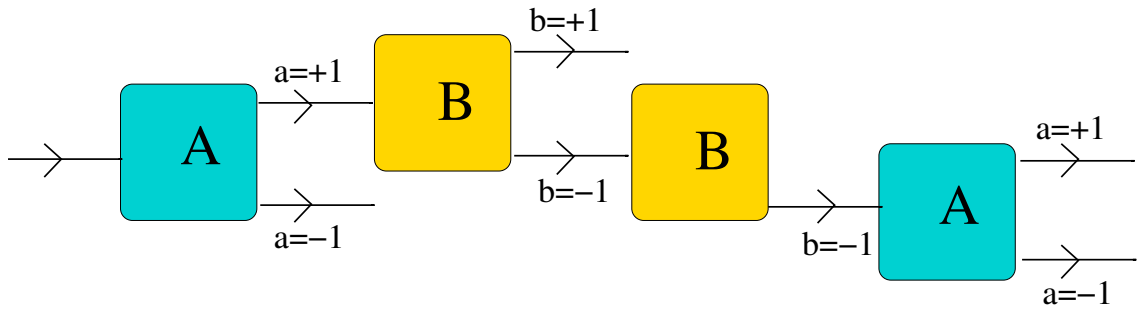
$$\hat{H} \longrightarrow \frac{\mu\hbar}{\sqrt{2}} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}.$$

Construct the representation of the propagator $\hat{U}(t, 0)$ in the same basis, and hence find the subsequent state vectors for:

(i) $|\psi(0)\rangle \longrightarrow (0, 1, 0)^\top$, (ii) $|\psi(0)\rangle \longrightarrow (1, 0, -1)^\top/\sqrt{2}$.

The eigenvectors of the matrix representation of \hat{H} were obtained in Q4(i) of Examples 1. Show that your results can also be obtained by first decomposing $|\psi(0)\rangle$ in the eigenbasis of \hat{H} .

5.



The picture above represents a series of measurements of observables associated with operators \hat{A} and \hat{B} , each of which has only two eigenvalues, ± 1 . In the basis $\{|a+\rangle, |a-\rangle\}$ the operators are given by

$$\hat{A} \xrightarrow{a} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad \hat{B} \xrightarrow{a} \begin{pmatrix} \cos 2\theta & \sin 2\theta \\ \sin 2\theta & -\cos 2\theta \end{pmatrix}$$

where θ is some real parameter. At each step, give the fraction of particles for which the measurement yields 1 or -1 .

[Hint: first show that the eigenvectors of \hat{B} can be represented as follows:

$$|b+\rangle \xrightarrow{a} \begin{pmatrix} \cos \theta \\ \sin \theta \end{pmatrix}, \quad |b-\rangle \xrightarrow{a} \begin{pmatrix} \sin \theta \\ -\cos \theta \end{pmatrix};$$

see also Q3(iv) of Examples 1.]

6. Use Ehrenfest's theorem to show that for the one-dimensional harmonic oscillator in any state (not necessarily a stationary state),

$$m \frac{d^2 \langle \hat{x} \rangle}{dt^2} = -k_s \langle \hat{x} \rangle,$$

where k_s is the spring constant. Comment on this result.

7. Consider a system with Hamiltonian $\hat{H} = \hbar\gamma \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$, and an observable $\hat{\Omega} = \frac{\hbar}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$.

Check that the state $\begin{pmatrix} \cos \gamma t \\ -i \sin \gamma t \end{pmatrix}$ is a solution of the time-dependent Schrödinger equation, and verify that Ehrenfest's theorem holds for the expectation value of $\hat{\Omega}$.

8. A particle of mass m is bound by a potential $V(\mathbf{r})$. Show that if the potential is spherically symmetric, $V = V(r)$, the expectation value of the angular momentum is conserved.

[You may use the commutation relations from Q8 of Examples 1.]

Now consider a non-spherically-symmetric potential $V = V_0(r) + zV_1(r)$. Show that $\langle \hat{L}_z \rangle$ is still conserved, and that

$$\frac{d}{dt} \langle \hat{L}_x \rangle = -\langle \hat{y} \hat{V}_1 \rangle \quad \text{and} \quad \frac{d}{dt} \langle \hat{L}_y \rangle = \langle \hat{x} \hat{V}_1 \rangle$$

Check that the right-hand sides correspond to the expectation values of the components of the torque, as expected.

9. This question uses the raising and lowering operators for the harmonic oscillator, \hat{a}^\dagger and \hat{a} .
- Verify the numerical coefficients in the expressions $\hat{a}^\dagger|n\rangle = \sqrt{n+1}|n+1\rangle$ and $\hat{a}|n\rangle = \sqrt{n}|n-1\rangle$, and show that these imply $\langle n|\hat{a} = \sqrt{n+1}\langle n+1|$ and $\langle n|\hat{a}^\dagger = \sqrt{n}\langle n-1|$.
 - Find the matrix elements $\langle m|\hat{x}|n\rangle$, $\langle m|\hat{p}|n\rangle$, $\langle m|\hat{x}^2|n\rangle$, $\langle m|\hat{p}^2|n\rangle$.
 - From the results of the previous part, find the uncertainty product $\Delta x\Delta p$ for a particle in the n th state, and comment on the result.
10. Verify that the definition

$$H_n(x) = \exp(x^2/2) \left(x - \frac{d}{dx} \right)^n \exp(-x^2/2)$$

does indeed generate the first few Hermite polynomials, as given in Q6 of Examples 1.

11. A state $|\lambda\rangle$ is an eigenstate of \hat{a} : $\hat{a}|\lambda\rangle = \lambda|\lambda\rangle$, for some complex λ . Find Δx and Δp .
[You may use the following as a check on your results: $\langle \hat{x} \rangle = \sqrt{2}x_0\text{Re}[\lambda]$.]
 Find an expression for the state $|\lambda\rangle$.
[Hint: writing $|\lambda\rangle = \sum_{n=0}^{\infty} c_n|n\rangle$, find a recurrence relation between the coefficients c_n , namely: $c_{n+1} = \lambda c_n / \sqrt{n+1}$. Don't worry about normalisation initially, but leave the first constant, c_0 , to be determined at the end.]

12. Consider the symmetric two-dimensional harmonic oscillator with potential $\frac{1}{2}m\omega^2(x^2 + y^2)$. Its energy eigenstates can be written $|n_x, n_y\rangle$ with energy $\hbar\omega(n_x + n_y + 1)$, n_x, n_y being non-negative integers.

What is the degeneracy of the state with energy $N\hbar\omega$, for positive integer N ?

Show that the Hermitian operator $\hat{L} = \hat{x}\hat{p}_y - \hat{y}\hat{p}_x$ can be written as

$$\hat{L} = i\hbar(\hat{a}_y^\dagger\hat{a}_x - \hat{a}_x^\dagger\hat{a}_y),$$

and that it commutes with the Hamiltonian.

In the subspace of states with energy $3\hbar\omega$, $\{|2, 0\rangle, |1, 1\rangle, |0, 2\rangle\}$, show that

$$\hat{L} \xrightarrow{N=3} \sqrt{2}\hbar \begin{pmatrix} 0 & -i & 0 \\ i & 0 & -i \\ 0 & i & 0 \end{pmatrix}.$$

Verify that the eigenvectors of this matrix are

$$\frac{1}{2} \begin{pmatrix} 1 \\ \sqrt{2}i \\ -1 \end{pmatrix}, \quad \frac{1}{2} \begin{pmatrix} 1 \\ -\sqrt{2}i \\ -1 \end{pmatrix} \quad \text{and} \quad \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ 1 \end{pmatrix}$$

and find the corresponding eigenvalues. Write down the position-space representations of these eigenstates, first in terms of Cartesian coordinates (x, y) and then in polars (r, ϕ) , where $x = r \cos \phi$ and $y = r \sin \phi$. Comment on your results.

13. For the symmetric three-dimensional harmonic oscillator, give expressions for the allowed energies and their degeneracies. Use them to predict the first three magic numbers in nuclei.

