

Lecture 17

Electron orbiting nucleus feels magnetic field $\mathbf{B} \propto \mathbf{L}$

→ spin-orbit coupling

$$\hat{H}_{\text{so}} = f(r) \hat{\mathbf{L}} \cdot \hat{\mathbf{S}}$$

In terms of squares of angular momentum operators

$$\hat{\mathbf{L}} \cdot \hat{\mathbf{S}} = \frac{1}{2} \left[\hat{\mathbf{J}}^2 - \hat{\mathbf{L}}^2 - \hat{\mathbf{S}}^2 \right] \quad \text{where} \quad \hat{\mathbf{J}} = \hat{\mathbf{L}} + \hat{\mathbf{S}}$$

Eigenstates of $\hat{\mathbf{J}}^2$, \hat{J}_z are eigenstates of $\hat{\mathbf{L}} \cdot \hat{\mathbf{S}}$

$$\hat{\mathbf{L}} \cdot \hat{\mathbf{S}} \psi_{JM_J} = \frac{\hbar^2}{2} \left[J(J+1) - L(L+1) - S(S+1) \right] \psi_{JM_J}$$

→ not mixed by \hat{H}_{so} (unlike degenerate eigenstates of \hat{L}_z , \hat{S}_z)

Using first-order perturbation theory

$$E_J = \mathcal{E} \frac{1}{2} \left[J(J+1) - L(L+1) - S(S+1) \right]$$

where $\mathcal{E} = \langle f(r) \rangle \hbar^2 > 0$ and so smaller J has lower energy