

**PHYS20352: Thermal and Statistical Physics**  
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**Example Sheet 3**

1. A quantity,  $n$  moles, of an ideal gas is taken through the following three stages of a reversible cycle:

- $a \rightarrow b$ : a quasistatic isochoric (i.e., at constant volume,  $V_a$ ) increase in pressure from  $P_a$  to  $P_b$ , with a corresponding increase in temperature, brought about by the absorption of an amount of heat  $Q_H$  from a series of external reservoirs ranging in temperature from  $T_a$  to  $T_b$ ;
- $b \rightarrow c$ : a quasistatic adiabatic expansion from volume  $V_a$  to  $V_c$ , involving a corresponding drop in temperature from  $T_b$  to  $T_c$ ; and finally
- $c \rightarrow a$ : a quasistatic isobaric (i.e., at constant pressure,  $P_a$ ) decrease in volume from  $V_c$  to  $V_a$  with a corresponding decrease in temperature, brought about by rejection of an amount of heat  $Q_C$  to a series of external reservoirs ranging in temperature from  $T_c$  to  $T_a$ .

- (a) Draw the cycle on a standard  $P$ - $V$  diagram.
- (b) Calculate the work done **by** the system,  $w$ , and the heat absorbed or rejected, during each of the three stages. [Note: You may find it useful to refer to the results of Sheet 2 Q2 for the quasistatic adiabatic process.]
- (c) Verify that the total work  $w$  done by the system is equal to the net heat absorbed, namely,  $(Q_H - Q_C)$ , or more specifically,

$$w = C_V(T_b - T_a) - C_P(T_c - T_a).$$

[Note: You may find it useful to refer to the results of Sheet 2 Q1(a).]

- (d) Hence show the efficiency,  $\eta \equiv w/Q_H$ , of the system as a heat engine is given by

$$\eta = 1 - \frac{\gamma(1-r)}{r(1-r^\gamma)} \frac{T_a}{T_b},$$

where  $\gamma \equiv C_P/C_V$ , and  $r = V_a/V_c$  is the inverse of the compression ratio.

2. (a) Use the first law of thermodynamics for a general gas (i.e., not necessarily an ideal gas) to derive the two following equations:

$$dQ = C_V dT + \left[ \left( \frac{\partial E}{\partial V} \right)_T + P \right] dV,$$

where  $C_V$  is the heat capacity at constant volume and

$$C_P = C_V + \left[ \left( \frac{\partial E}{\partial V} \right)_T + P \right] \left( \frac{\partial V}{\partial T} \right)_P,$$

where  $C_P$  is the heat capacity at constant pressure.

(b) Now consider one mole of a van der Waals gas, obeying the equation of state

$$\left( P + \frac{a}{V^2} \right) (V - b) = RT.$$

Its internal energy is given by

$$E = cT - \frac{a}{V},$$

where  $a, b$ , and  $c$  are constants, and  $R$  is the gas constant. Hence show that the molar heat capacities at constant volume and constant pressure are given respectively, by

$$C_V = c, \quad C_P = c + \frac{R}{1 - 2a(V - b)^2/(RTV^3)}.$$