## PHYS30201 Advanced Quantum Mechanics: Examples 1 - Revision

Most of these questions are revision; you will find sections A.1-A. 4 of the web notes useful for these.

1. (a) Prove that the eigenvalues $\left\{\omega_{i}\right\}$ of a Hermitian operator $\widehat{\Omega}$ are real and the (normalised, non-degenerate) eigenvectors $\left\{\left|\omega_{i}\right\rangle\right\}$ are orthogonal.
(b) If $|\alpha\rangle=\sum_{i} a_{i}\left|\omega_{i}\right\rangle$ and $|\beta\rangle=\sum_{i} b_{i}\left|\omega_{i}\right\rangle$, show that $\langle\beta \mid \alpha\rangle=\sum_{i} b_{i}^{*} a_{i}=(\langle\alpha \mid \beta\rangle)^{*}$.
(c) Prove also that $\sum_{i}\left|\omega_{i}\right\rangle\left\langle\omega_{i}\right|=\widehat{I}$. (Hint - show that acting on an arbitrary ket $|\alpha\rangle$ gives $|\alpha\rangle$ again.)
(d) Another Hermitian operator $\widehat{\Theta}$ has matrix elements $\left\langle\omega_{i}\right| \widehat{\Theta}\left|\omega_{j}\right\rangle=\theta_{i j}$. Show that $\langle\beta \mid \widehat{\Theta} \alpha\rangle=$ $\sum_{i j} \theta_{i j} b_{i}^{*} a_{j}=\langle\widehat{\Theta} \beta \mid \alpha\rangle$, and also that $\sum_{i j}\left|\omega_{i}\right\rangle \theta_{i j}\left\langle\omega_{j}\right| \equiv \widehat{\Theta}$.
2. A particle is in a state $|\psi\rangle$ when the observable corresponding to $\widehat{\Omega}$ is measured. Show that the following two statements are equivalent (i.e. assume one and show that the other is a consequence of it):
(I) the probability of getting a result $\omega_{i}$ is $\left|\left\langle\omega_{i} \mid \psi\right\rangle\right|^{2}$;
(II) the expectation value (ensemble average) is $\langle\psi| \widehat{\Omega}|\psi\rangle$.
3. (a) Verify that $\langle\alpha \mid \beta\rangle=\int \alpha\left(x^{\prime}\right)^{*} \beta\left(x^{\prime}\right) \mathrm{d} x^{\prime}$. (Hint - use the appropriate identity operator.)
(b) Verify that $\langle x| \hat{p}\left|x^{\prime}\right\rangle=i \hbar \frac{\mathrm{~d}}{\mathrm{~d} x^{\prime}} \delta\left(x-x^{\prime}\right)$ implies $\langle x| \hat{p}|\alpha\rangle=-i \hbar \frac{\mathrm{~d}}{\mathrm{~d} x} \alpha(x)$. (Hint - as above!)
(c) Verify using the previous result that $\langle x|[\hat{x}, \hat{p}]|\alpha\rangle=i \hbar \alpha(x)$ (ie not just by replacing the commutator by $i \hbar)$.
(d) Verify that if $\tilde{\alpha}(p)$ is the Fourier transform of $\alpha(x), p \tilde{\alpha}(p)$ is the Fourier transform of $-i \hbar \frac{\mathrm{~d}}{\mathrm{~d} x} \alpha(x)$. (Use the convention that has $-i k x$ in the exponent for the transform, and $i k x$ for the inverse.)
(e) If $\phi(x)=\left(\pi x_{0}^{2}\right)^{-1 / 4} \exp \left(-x^{2} /\left(2 x_{0}^{2}\right)\right)$, find $\tilde{\phi}(p)$ and check the normalisation.
4. A hydrogen atom is prepared in an initial state described by the wave function

$$
\psi_{I}(\mathbf{r})=\frac{1}{\sqrt{96 \pi a_{0}^{5}}} r \exp \left(-\frac{r}{2 a_{0}}\right)
$$

A measurement is made of the energy. What is the probability of obtaining the ground state energy of -13.6 eV ? What other values do you think might be obtained? What would the probability be if instead the initial state were $\sqrt{3} \cos (\theta) \psi_{I}(\mathbf{r})$ ? (Hint: You will need to know the ground-state wave function of hydrogen to answer this. The first answer works out as 0.234 .)
5. A particle has $\operatorname{spin} s=\frac{3}{2}$ and orbital angular momentum $l=1$.
(a) List all the allowed values of $\left\{m_{s}, m_{l}\right\}$. List also all the allowed values of $\left\{j, m_{j}\right\}$, where the vector operator $\widehat{\mathbf{J}}$ is the total angular momentum operator of the system. Check that the two lists are equal in length, and also that there are the same number of pairs with a given $m_{j}$ in the second list as with a given $m_{s}+m_{l}$ in the first.
(b) We can use either set of quantum numbers to define a basis. In the first case we denote the basis states by $\left|s m_{s}\right\rangle \otimes\left|l m_{l}\right\rangle$, and in the second case by $\left|l, s ; j m_{s}\right\rangle$. By looking up the relevant Clebsch-Gordan coefficients, write the state $\left|1, \frac{3}{2} ; \frac{1}{2} \frac{1}{2}\right\rangle$ as a sum of states from the first basis.
(c) The matrix representations of the angular momentum operators in the basis $\{|11\rangle,|10\rangle$, $|1-1\rangle\}$ are

$$
L_{x}=\frac{\hbar}{\sqrt{2}}\left(\begin{array}{ccc}
0 & 1 & 0 \\
1 & 0 & 1 \\
0 & 1 & 0
\end{array}\right) \quad L_{y}=\frac{\hbar}{\sqrt{2}}\left(\begin{array}{ccc}
0 & -i & 0 \\
i & 0 & -i \\
0 & i & 0
\end{array}\right) \quad L_{z}=\hbar\left(\begin{array}{ccc}
1 & 0 & 0 \\
0 & 0 & 0 \\
0 & 0 & -1
\end{array}\right)
$$

Use these to show that the state

$$
\sqrt{\frac{2}{3}}|11\rangle \otimes\left|\frac{1}{2}-\frac{1}{2}\right\rangle-\sqrt{\frac{1}{3}}|10\rangle \otimes\left|\frac{1}{2} \frac{1}{2}\right\rangle
$$

is an eigenfunction of $\widehat{\mathbf{J}}^{2}$ with eigenvalue $\frac{3}{4} \hbar^{2}$.
6. The states $\{|n\rangle\}$ are the eigenstates of a harmonic oscillator of mass $m$ and oscillator frequency $\omega$. Where appropriate use the length $x_{0}=\sqrt{\hbar /(m \omega)}$ in your answers.
(a) Use creation and annihilation operators to find the matrix elements $\langle m| \hat{x}|n\rangle,\langle m| \hat{p}|n\rangle$, $\langle m| \hat{x}^{2}|n\rangle,\langle m| \hat{p}^{2}|n\rangle$.
(b) From the results of the previous part, find the uncertainty product $\Delta x \Delta p$ for a particle in the $n$th state, and comment on the result.
7. This question is a little involved, but we will use the exact results obtained here to check our answers when we treat the same problem in perturbation theory later on.
Consider the new Hamiltonian corresponding to adding a linear potential to a harmonic oscillator, so that $\widehat{H}^{\prime}=\frac{\hat{p}^{2}}{2 m}+\frac{1}{2} m \omega^{2} \hat{x}^{2}-\lambda \hat{x}$.
(a) Rewrite the Hamiltonian to make clear that it is simply a new harmonic potential, and find expressions for the shift in the origin $\delta$ and the energy offset $\Delta E$.
(b) Show that the relationship between the creation and annihilation operators for the new system, $\hat{b}^{\dagger}$ and $\hat{b}$, and those for the original system $\hat{a}^{\dagger}$ and $\hat{a}$, is $\hat{b}^{\dagger}=\hat{a}^{\dagger}-\delta /\left(\sqrt{2} x_{0}\right)$ and $\hat{b}=\hat{a}-\delta /\left(\sqrt{2} x_{0}\right)$.. Use these to find the coefficients in the expansion of the new ground state $\left|0^{\prime}\right\rangle$, annihilated by $\hat{b}$, in terms of the states of the original system $|n\rangle$, and check that the sum of their magnitudes squared is unity. (Hint: You will need to use the wave function to calculate the overlap of the old and new ground states $\left\langle 0 \mid 0^{\prime}\right\rangle$, but all the other overlaps $\langle n \mid 0\rangle$ can be reduced to this one with operator methods.)

